

# Towards Three-Dimensional Conformal Probability

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Conference in honor of Vincent Rivasseau  
Paris, November 25, 2015

- ① The Brydges-Mitter-Scoppola (BMS) model in 3D
- ② The  $p$ -adic hierarchical model

# 1) Heuristic definition of the BMS model

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$$\frac{1}{Z} \exp \left( -\frac{1}{2} \langle \phi, (-\Delta)^\alpha \phi \rangle - \int_{\mathbb{R}^3} \{g\phi(x)^4 + \mu\phi(x)^2\} d^3x \right) D\phi$$

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**Main focus:** the scale invariant theory corresponding to a nontrivial RG fixed point (BMS in CMP 2003).

Should correspond to critical scaling limit of **long-range** three-dimensional Ising model with ferromagnetic interactions  $J_{\mathbf{x},\mathbf{y}} \sim |\mathbf{x} - \mathbf{y}|^{-(d+\sigma)}$ ,  $d = 3$ ,  $\sigma = \frac{3+\epsilon}{2}$ .

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Fix rescaling ratio  $L > 1$ , integer.

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Given **bare ansatz**  $(g_r, \mu_r)_{r \in \mathbb{Z}}$ , one has well defined probability measures  $d\nu_{r,s}(\phi)$  whose Radon-Nikodym derivative with respect to  $d\mu_{C_r}(\phi)$  is proportional to

$$\exp \left( - \int_{\mathbb{R}^3} \rho_{IR,s}(x) \{ g_r : \phi^4 : (x) + \mu_r : \phi^2 : (x) \} d^3x \right)$$

with Wick ordering relative to  $\mu_{C_r}$ .

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The scale invariant BMS measure should be the weak limit  $\nu_\phi = \lim_{r \rightarrow -\infty} \lim_{s \rightarrow \infty} \nu_{r,s}$  for a well chosen bare ansatz which mimics the scaling limit of a critical theory on the unit lattice.



## Conjecture 1:

Set  $[\phi] = \frac{3-\epsilon}{4}$  for  $\epsilon$  positive and small.

Then there exists a nonempty interval  $I \subset (0, \infty)$  and a function  $\mu_c : I \rightarrow \mathbb{R}$  such that for all  $g \in I$ , if one sets  $g_r = L^{-r(3-4[\phi])}g$  and  $\mu_r = L^{-r(3-2[\phi])}\mu_c(g)$ , the weak limit  $\nu_\phi$  exists and is non-Gaussian, translation-invariant,  $O(3)$ -invariant, OS positive, and scale-invariant with exponent  $[\phi]$ , i.e.,  $\lambda^{[\phi]}\phi(\lambda \cdot) \stackrel{dd}{=} \phi(\cdot)$  for all  $\lambda > 0$ .

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Measure constructed on finite torus by Mitter ( $\sim 2004$ ) using fixed point obtained by BMS, CMP 2003.

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The moments

$$S_n(f_1, \dots, f_n) = \langle \phi(f_1) \cdots \phi(f_n) \rangle = \int_{S'(\mathbb{R}^3)} \phi(f_1) \cdots \phi(f_n) d\mu(\phi)$$

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By the nuclear theorem  $S_n$  can be seen as an element of  $S'(\mathbb{R}^{3n})$ .

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Conjecture 2:

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Conj. 3 is a precise formulation of predictions in “Conformal invariance in the long-range Ising model” by Paulos, Rychkov, van Rees and Zan, arXiv:1509.00008[hep-th] — >

**Higher-dimensional conformal bootstrap.**



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**One-to-one correspondence:**  $f \in \mathcal{M}(\mathbb{R}^3) \leftrightarrow$  hyperbolic isometry of the bulk  $\mathbb{B}^4$  or  $\mathbb{H}^4$ .

- ① The Brydges-Mitter-Scoppola (BMS) model in 3D
- ② The  $p$ -adic hierarchical model



# 1) Hierarchical continuum

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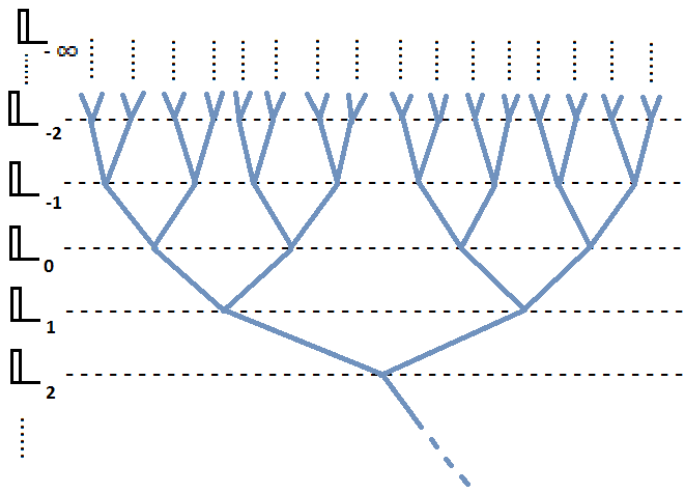
Let  $\mathbb{L}_k$ ,  $k \in \mathbb{Z}$ , be the set of boxes  $\prod_{i=1}^d [a_i p^k, (a_i + 1)p^k)$  for  $a_1, \dots, a_d \in \mathbb{N}$ . The cubes in  $\mathbb{L}_k$  form a partition of the octant  $[0, \infty)^d$ .

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Then  $\mathbb{T} = \cup_{k \in \mathbb{Z}} \mathbb{L}_k$  naturally has the structure of a **doubly** infinite tree organized in layers or generations  $\mathbb{L}_k$ :



Picture for  $d = 1$ ,  $p = 2$

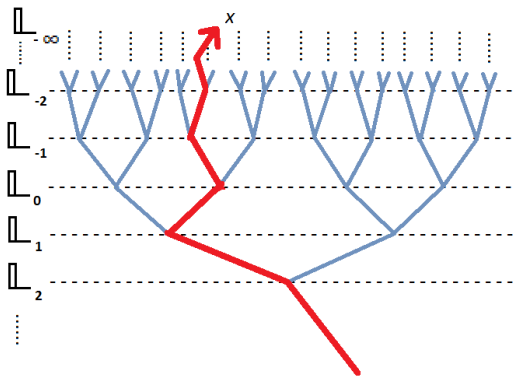
Now **forget about**  $[0, \infty)^d$  and  $\mathbb{R}^d$ .

Define the substitute for the continuum  $\mathbb{Q}_p^d :=$  set of leafs at infinity " $\mathbb{L}_{-\infty}$ ".

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More precisely, this is the set of upward paths in the tree.



A path representing some  $x \in \mathbb{Q}_p^d$

A point  $x \in \mathbb{Q}_p^d$  encoded by sequence  $(a_n)_{n \in \mathbb{Z}}$ ,  
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### Caution! perverse notation ahead

$a_n$  represents local coordinates of  $\mathbb{L}_{-n-1}$  box inside  $\mathbb{L}_{-n}$  box.

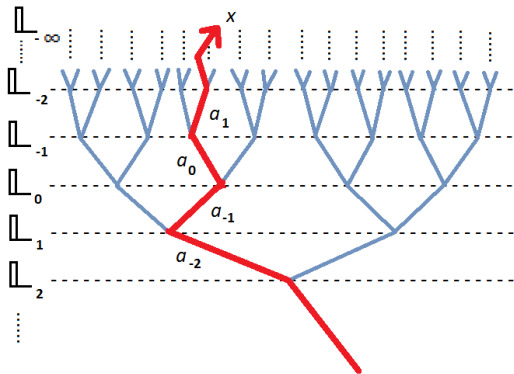


A point  $x \in \mathbb{Q}_p^d$  encoded by sequence  $(a_n)_{n \in \mathbb{Z}}$ ,  
 $a_n \in \{0, 1, \dots, p-1\}^d$ .

Let  $0 \in \mathbb{Q}_p^d$  correspond to sequence with all digits equal to zero.

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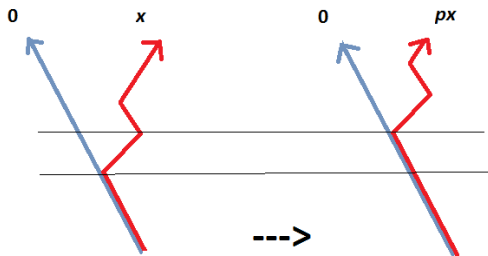


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if  $x = (a_n)_{n \in \mathbb{Z}}$  then  $px = (a_{n-1})_{n \in \mathbb{Z}}$ , i.e., **upward** shift.

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Likewise  $p^{-1}x$  is downward shift and so on for defining  $p^k x$ ,  
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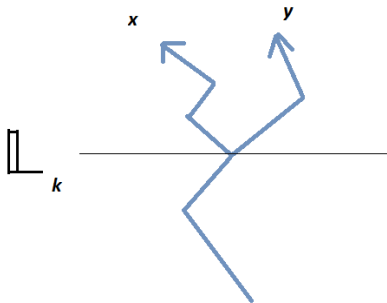
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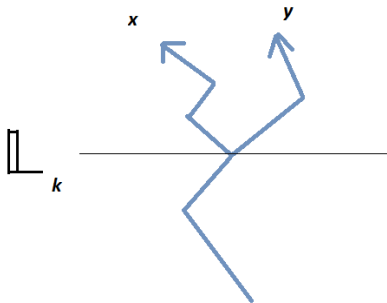
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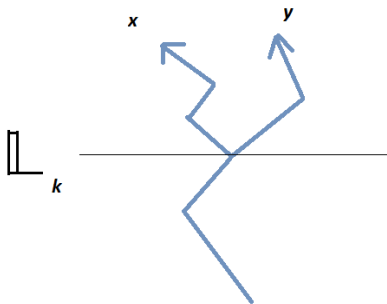
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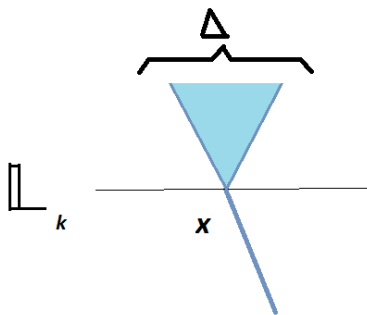
also define  $|x| := |x - 0|$ . Because of the strange notation

$$|px| = p^{-1}|x|$$



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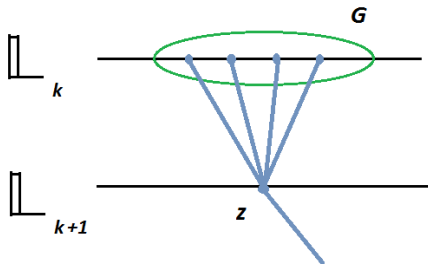
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Construction: take product of uniform probability measures on  $(\{0, 1, \dots, p-1\}^d)^\mathbb{N}$  for  $\overline{B}(0, 1)$  and similarly for other balls of radius 1, then collate.

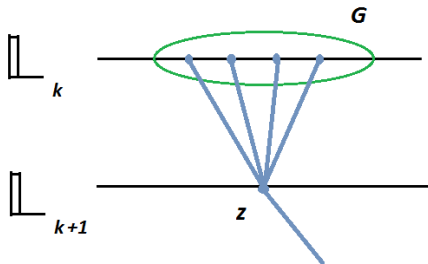
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Massless Gaussian field  $\phi(\mathbf{x})$ ,  $\mathbf{x} \in \mathbb{Q}_p^d$  with engineering scaling dimension  $[\phi]$  is

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only formal since  $\phi$  not defined pointwise. **Need random distributions.**

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where for  $t_- \leq t_+$ ,  $S_{t_-, t_+}(\mathbb{Q}_p^d)$  is space of functions which are constant in closed boxes of radius  $p^{t_-}$  and support in  $\overline{B}(0, p^{t_+})$ .

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Topology generated by the set of **all** seminorms.



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Probability theory on  $S'(\mathbb{Q}_p^d)$  is very nice!

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- ⑥  $S'(\mathbb{Q}_p^d) \times S'(\mathbb{Q}_p^d) \simeq S'(\mathbb{Q}_p^d)$  so same tools work for joint law of pair of distributional random fields, e.g.,  $(\phi, N[\phi^2])$ .

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If law of  $\phi(\cdot)$  is  $\mu_{C_0}$ , then law of  $L^{-r[\phi]}\phi(L^r \cdot)$  is  $\mu_{C_r}$ .

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Let

$$V_{r,s}(\phi) = \int_{\Lambda_s} \{g_r : \phi^4 :_{C_r}(x) + \mu_r : \phi^2 :_{C_r}(x)\} d^3x$$

and define the probability measure

$$d\nu_{r,s}(\phi) = \frac{1}{Z_{r,s}} e^{-V_{r,s}(\phi)} d\mu_{C_r}(\phi)$$

Let  $\phi_{r,s}$  random variable in  $S'(\mathbb{Q}_p^3)$  sampled according to  $\nu_{r,s}$  and define **square** field  $N_r[\phi_{r,s}^2]$  which is **deterministic**  $S'(\mathbb{Q}_p^3)$ -valued function of  $\phi_{r,s}$  given by

$$N_r[\phi_{r,s}^2](j) = Z_2^r \int_{\mathbb{Q}_p^3} \{ Y_2 : \phi_{r,s}^2 :_{C_r}(x) - Y_0 L^{-2r[\phi]} \} j(x) d^3x$$

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Our main result concerns the limit law of the pair  $(\phi_{r,s}, N_r[\phi_{r,s}^2])$  in  $S'(\mathbb{Q}_p^3) \times S'(\mathbb{Q}_p^3)$  when  $r \rightarrow -\infty, s \rightarrow \infty$  **regardless** of the order of limits.

Will need approximate fixed point coupling

$$\bar{g}_* = \frac{p^\epsilon - 1}{36L^\epsilon(1 - p^{-3})}$$

## 8) Results

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Mixed correlations satisfy in sense of distributions

$$\begin{aligned} & \langle \phi(L^{-k}x_1) \cdots \phi(L^{-k}x_n) N[\phi^2](L^{-k}y_1) \cdots N[\phi^2](L^{-k}y_m) \rangle \\ &= L^{-(n[\phi]+m[\phi^2])k} \langle \phi(x_1) \cdots \phi(x_n) N[\phi^2](y_1) \cdots N[\phi^2](y_m) \rangle \end{aligned}$$

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$[\phi^2] - 2[\phi] = \frac{1}{3}\epsilon + o(\epsilon)$  as expected in Euclidean BMS model.



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The law  $\nu_{\phi \times \phi^2}$  of  $(\phi, N[\phi^2])$  is independent of  $g$ : **universality**

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$\nu_{\phi \times \phi^2}$  is **fully** scale invariant, i.e., invariant under action of  $p^{\mathbb{Z}}$  instead of just  $L^{\mathbb{Z}}$ . Moreover,  $\mu(g)$  and  $[\phi^2]$  independent of RG step  $L$ .

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Two point functions given as distributions by

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$$\langle \phi(x)\phi(y) \rangle = \frac{c_1}{|x-y|^{2[\phi]}}$$

$$\langle N[\phi^2](x) N[\phi^2](y) \rangle = \frac{c_2}{|x-y|^{2[\phi^2]}}$$

Note that  $3 - 2[\phi^2] = 3 - \frac{1}{3}\epsilon + o(\epsilon) \rightarrow$  still  $L^{1,loc}$  !

### Theorem 3: A.A., May 2015

Let  $\psi_i$  denote  $\phi$  or  $N[\phi^2]$ . Then for every mixed correlation  $\exists$  smooth function  $\langle \psi_1(z_1) \cdots \psi_n(z_n) \rangle$  on  $(\mathbb{Q}_p^3)^n \setminus \text{Diag}$  which is locally integrable (even on  $\text{Diag}$ ) such that

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for all test functions  $f_1, \dots, f_n \in \mathcal{S}(\mathbb{Q}_p^3)$ .

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Namely,  $\nu_{\phi \times \phi^2}$  satisfies the DPC property or the analogue of Conj. 2.

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$p$ -adic Möbius group: generated by isometries (for the ultrametric  $|x - y|$ ,  $x, y \in \mathbb{Q}_p^3$ ), scaling transformations  $x \mapsto p^k x$ ,  $k \in \mathbb{Z}$ , and the unit-sphere inversion  $J(x) = |x|^2 x$ .

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Alternatively, one can define the absolute cross-ratio using the ultrametric, then  $\mathcal{M}(\mathbb{Q}_p^3)$  is the group of bijections of  $\widehat{\mathbb{Q}_p^3} = \mathbb{Q}_p^3 \cup \{\infty\}$  which preserve the cross-ratio.

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The AdS bulk is the tree  $\mathbb{T}$  with the graph distance. Analogue of hyperbolic metric.

## Mumford-Manin-Drinfeld Cross-Ratio Lemma

$$CR(x_1, x_2, x_3, x_4) := \frac{|x_1 - x_3| |x_2 - x_4|}{|x_1 - x_4| |x_2 - x_3|} = p^{-\delta(x_1 \rightarrow x_2; x_3 \rightarrow x_4)},$$

where  $\delta(x_1 \rightarrow x_2; x_3 \rightarrow x_4)$  is the number of common edges in bi-infinite paths  $x_1 \rightarrow x_2$  and  $x_3 \rightarrow x_4$ , counted positively if orientations agree and negatively otherwise.

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Using the cross-ratio lemma, one can again establish a one-to-one correspondence:  $f \in \mathcal{M}(\mathbb{Q}_p^3) \leftrightarrow$  hyperbolic isometry of the bulk  $\mathbb{T}$ .

## Mumford-Manin-Drinfeld Cross-Ratio Lemma

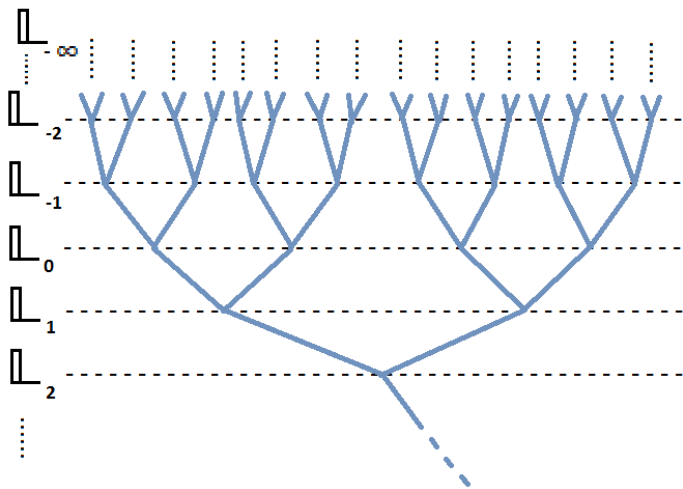
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Rigorous RG for space-dependent couplings in ACG 2013 –  $\rightarrow$  space-dependent cut-offs –  $\rightarrow$  Conj. 3. by showing equivalence of usual upper half-space cut-off with conformal ball cut-off.





The tree again

## 10) The last slide

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Cher Vincent,

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Je te souhaite encore beaucoup d'années,

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Je te souhaite encore beaucoup d'années,  
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Joyeux Anniversaire!